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THE EXACT DISTRIBUTION OF EXOGENOUS

VARIABLE COEFFICIENT ESTIMATORS

Peter C. B. Phillips

THE EXACT DISTRIBUTION OF EXOGENOUS VARIABLE COEFFICIENT ESTIMATORS

by

P. C. B. Phillips*

Cowles Foundation Yale University

O. ABSTRACT

This paper derives the exact probability density function of the instrumental variable (IV) estimator of the exogenous variable coefficient vector in a structural equation containing n+1 endogenous variables and N degrees of overidentification. A leading case of the general distribution that is more amenable to analysis and computation is also presented. Conventional classical assumptions of normally distributed errors and nonrandom exogenous variables are employed.

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1. INTRODUCTION

Substantial progress has been made in recent years on the exact distribution theory of econometric estimators and test statistics in simultaneous equations models. The latest results cover general specifications of single equation estimation which allow for the presence of any number of endogenous variables and an arbitrary degree of (either apparent or effective) equation overidentification. Thus, in earlier papers, the author (1980, 1983a, 1983b) has given the exact distributions of the instrumental variable (IV) and limited information maximum likelihood (LIML) estimators in this general setting; Rhodes (1981) extracted the exact density of the limited information identifiability test statistic; and Hillier, Kinal and Srivastava (1983) have provided exact moment formulae for the marginal distributions of the IV estimator. For a recent review of these and other developments in the field the reader is referred to Phillips (1983c).

The structural equation distribution theory cited above concentrates on the estimated coefficients of the endogenous variables. This is natural because these coefficients form the nucleus of the simultaneity problem and are therefore our primary concern. The coefficients of the exogenous variables are also an important, if subsidiary, component in the study of structural estimation. So far, our knowledge of the distribution of the estimated exogenous variable coefficients comes from the moment formulae that can be deduced from the equations that define these estimators in terms of the estimated endogenous variable coefficients (see, for instance, Phillips (1983c)). But the exogenous variable coefficients gain in significance in the transition from

structure to reduced form. And an understanding of the distribution of these estimated coefficients provides an important stepping stone to the study of the estimated reduced forms.

The present paper derives the exact probability density function (p.d.f.) of the estimated exogenous variable coefficients in a general single equation setting when the estimation method is instrumental variables. Conventional assumptions of normally distributed errors and nonrandom exogenous variables are employed. A leading case is presented in Section 3.

2. THE MODEL AND NOTATION

We work with the structural equation

(1)
$$y_1 = Y_2 \beta + Z_1 \gamma + u$$

where $y_1(T \times 1)$ and $Y_2(T \times n)$ are an observation vector and observation matrix, respectively, of n+1 included endogenous variables, Z_1 , is a $T \times K_1$ matrix of included exogenous variables and u is a random disturbance vector. The reduced form of (1) is written:

(2)
$$[y_1 : Y_2] = [z_1 : z_2]_{\pi_{21}}^{\pi_{11}} \xrightarrow{\Pi_{12}} + [v_1 : V_2] = Z\Pi + V$$

where \mathbf{Z}_2 is a $\mathbf{T} \times \mathbf{K}_2$ matrix of exogenous variables excluded from (1). The rows of the reduced form matrix V are assumed to be independent identically distributed normal random vectors. We assume that standardizing transformations have been carried out so that the covariance matrix of each row of V is the identity matrix and $\mathbf{T}^{-1}\mathbf{Z}^{\dagger}\mathbf{Z} = \mathbf{I}_{\mathbf{K}}$ where

 $K=K_1+K_2$. These transformations involve no loss of generality and their effect on the parameterization and resulting estimator distributions are fully discussed in Phillips (1983c). We assume that $K_2 \geq n$ and denote the degree of overidentification by $N=K_2-n$. Finally we note that the relationship between (1) and (2) and the implied restrictions on (1) yield the equations:

(3)
$$\pi_{11} - \pi_{12}\beta = \gamma$$
, $\pi_{21} - \pi_{22}\beta = 0$.

We define $H = [Z_1 \ \vdots \ Z_3]$, where Z_3 $(T \ X \ K_3)$ is a submatrix of Z_2 and $Z_3 \ge n$. The IV estimators of β and γ in (1) obtained by using H as the matrix of instruments are:

(4)
$$\beta_{IV} = (Y_2^{\dagger} Z_3 Z_3^{\dagger} Y_2)^{-1} (Y_2^{\dagger} Z_3 Z_3^{\dagger} Y_1),$$

(5)
$$\gamma_{IV} = T^{-1}Z_{1}^{\dagger}y_{1} - T^{-1}Z_{1}^{\dagger}Y_{2}\beta_{IV}.$$

The number of surplus instruments in this estimation is denoted by $L \,=\, K_{\,3}\, - n \ .$

THE LEADING CASE

We define the matrix variate

(6)
$$[b : B] = [T^{-1}Z_1^{\dagger}y_1 : T^{-1}Z_1^{\dagger}Y_2] = T^{-1}Z_1^{\dagger}X$$

which is normal with mean matrix $[\pi_{11}:\Pi_{12}]$ and covariance matrix $I_{K_1(n+1)}$. From equation (B3) in Phillips (1980) we know the p.d.f. of β_{IV} to be:

$$(7) \quad pdf(\mathbf{r}) = \frac{\text{etr}\left\{-\frac{T}{2}(\mathbf{I}+\beta\beta^{\dagger})\overline{\Pi}_{22}^{\dagger}\overline{\Pi}_{22}\right\}\Gamma_{n}\left(\frac{\mathbf{L}+\mathbf{n}+1}{2}\right)}{\pi^{n/2}(\mathbf{1}+\mathbf{r}^{\dagger}\mathbf{r})^{(\mathbf{L}+\mathbf{n}+1)/2}}$$

$$\cdot \sum_{\mathbf{j}=0}^{\infty} \sum_{\mathbf{k}=0}^{\infty} \sum_{\boldsymbol{\varphi} \in \mathbf{j} \cdot \kappa} \frac{\left(\frac{L}{2}\right)_{\mathbf{j}}\left(\frac{\mathbf{L}+\mathbf{n}+1}{2}\right)_{\kappa} \theta_{\boldsymbol{\varphi}}^{\mathbf{j},\kappa}}{\mathbf{j}!\mathbf{k}!\Gamma_{n}\left(\frac{\mathbf{L}+\mathbf{n}}{2}, \boldsymbol{\varphi}\right)}$$

$$\cdot C_{\boldsymbol{\varphi}}^{\mathbf{j},\kappa}\left(\frac{T_{\overline{\mathbf{l}}}}{2^{2}}2^{\beta\beta^{\dagger}}\overline{\Pi}_{22}, \frac{T_{\overline{\mathbf{l}}}}{2^{2}}(\mathbf{I}+\beta\mathbf{r}^{\dagger})(\mathbf{I}+\mathbf{r}\mathbf{r}^{\dagger})^{-1}(\mathbf{I}+\mathbf{r}\beta^{\dagger})\overline{\Pi}_{22}^{\dagger}\right).$$

In this expression $C_{\phi}^{j,\kappa}$ is an invariant polynomial in the elements of its two argument matrices. Such polynomials were introduced by Davis (1979, 1980) to extend the zonal polynomials and the reader is referred to his articles for a detailed presentation of their properties, together with a definition of the constants $\theta_{\phi}^{j,\kappa}$ that appear in (7). ϕ is a partition of the integer f = j+k into $\leq n$ parts, κ is a partition of k into $\leq n$ parts and the notation $\phi \in (j,\kappa)$ which is defined by Davis (1979) relates the two sets of partitions in the summation. The matrix $\overline{\pi}_{22}$ in (7) depends only on the submatrix $\overline{\pi}_{22}$ of reduced form coefficients; it is defined in Phillips (1980).

Since $Z_1^{\dagger}Z_3 = 0$, $Z_1^{\dagger}X$ is statistically independent of β_{IV} (which is a rational function of the elements of $Z_3^{\dagger}X$). The joint p.d.f. of (b,B,β_{IV}) is therefore given by:

(8)
$$pdf(b,B,r) = (2\pi)^{-K_1/2-K_1} n/2 \exp \left\{ -(b-\pi_{11})!(b-\pi_{11})/2 \right\}$$

$$\cdot etr \left\{ -\frac{1}{2}(B-\Pi_{12})!(B-\Pi_{12}) \right\} pdf(r) .$$

We examine first the leading case that is characterized by the null hypothesis

(9)
$$H_0: \Pi_{12} = 0$$
, $\Pi_{22} = 0$.

Under H_0 , the rank condition for identification of (1) fails, the parameter vector β is no longer identifiable and estimation by IV proceeds under conditions of only apparent overidentification. On the other hand, $\gamma = \pi_{11}$ under H_0 . Thus, the analysis of this leading case can provide insight into the effects of misspecification through the imposition of erroneous simultaneity. In particular, estimation of γ can be achieved by conventional least squares regression of the reduced form under H_0 or by instrumental variables upon (1) leading to $\gamma_{\rm IV}$. Comparison of the distributions of the two estimators will illustrate the effects of the misspecification.

Under H_0 , the density (7) reduces to:

(10)
$$pdf(r) = \frac{\Gamma\left(\frac{L+n+1}{2}\right)}{\pi^{n/2}\Gamma\left(\frac{L+1}{2}\right)(1+r!r)(L+n+1)/2}$$

and (8) becomes

(11)
$$pdf(b,B,r) = \frac{\Gamma(\frac{L+n+1}{2})exp\{-(b-\gamma)!(b-\gamma)/2\}}{\frac{K_1(n+1)/2}{\pi}\frac{K_1(n+1)/2+n/2}{\Gamma(\frac{L+1}{2})}} \frac{etr\{-\frac{1}{2}B!B\}}{(1+r!r)^{(L+n+1)/2}}$$

We write γ_{IV} as s = b - Br (compare (5)) and under this transformation we deduce:

(12)
$$pdf(s,B,r) = \frac{\Gamma\left(\frac{L+n+1}{2}\right)etr\left\{-\frac{1}{2}B^{\dagger}B\right\}}{\frac{K_{1}(n+1)/2}{\pi}\frac{K_{1}(n+1)/2+n/2}{\Gamma\left(\frac{L+1}{2}\right)}} \frac{exp\left\{-\frac{1}{2}(s+Br-\gamma)^{\dagger}(s+Br-\gamma)\right\}}{(1+r^{\dagger}r)^{(L+n+1)/2}}$$

$$= \frac{\Gamma\left(\frac{L+n+1}{2}\right)exp\left\{-\frac{1}{2}(s-\gamma)^{\dagger}(s-\gamma)\right\}}{\frac{K_{1}(n+1)/2}{\pi}\frac{K_{1}(n+1)/2+n/2}{\Gamma\left(\frac{L+1}{2}\right)(1+r^{\dagger}r)^{(L+n+1)/2}}}$$

$$\cdot etr\left\{-\frac{1}{2}(I+rr^{\dagger})\left[B-(s-\gamma)r^{\dagger}(I+rr^{\dagger})^{-1}\right]^{\dagger}\left[B-(s-\gamma)r^{\dagger}(I+rr^{\dagger})^{-1}\right]\right\}$$

$$\cdot exp\left\{\frac{1}{2}(s-\gamma)^{\dagger}(s-\gamma)r^{\dagger}(I+rr^{\dagger})^{-1}r\right\}.$$

Integrating out B we find

$$pdf(s,r) = \frac{\Gamma\left(\frac{L+n+1}{2}\right)exp\left\{-\frac{1}{2}(s-\gamma)'(s-\gamma)\right\}}{\frac{K_1/2}{\pi}\frac{K_1/2+n/2}{\Gamma\left(\frac{L+1}{2}\right)(1+r'r)}\frac{exp\left\{\frac{1}{2}(s-\gamma)'(s-\gamma)r'r/(1+r'r)\right\}}{\left[\det(I+\tau\tau')\right]}}{\left[\det(I+\tau\tau')\right]}$$

(13)
$$= \frac{\Gamma\left(\frac{L+n+1}{2}\right) \exp\left\{-\frac{1}{2}(s-\gamma)'(s-\gamma)\right\}}{(2\pi)^{K_1/2} \pi^{n/2} \Gamma\left(\frac{L+1}{2}\right) (1+r'r)^{(L+n+1+K_1)/2}} \sum_{j=0}^{\infty} \frac{\left\{\frac{1}{2}(s-\gamma)'(s-\gamma)\right\}^{j} (r'r)^{j}}{j! (1+r'r)^{j}} .$$

We transform $r \to (m,h)$ according to the decomposition $r = (r^*r)^{1/2}(r/(r^*r)^{1/2})$ = $m^{1/2}h$. The measure changes in accordance with the relation

$$dr = 2^{-1}m^{(n-2/2)}dm(dh)$$

where (dh) denotes the invariant measure over the Stiefel manifold $V_{1,n}$ (see James (1954), equation (8.19)). With this transformation we integrate out (m,h) as follows:

$$(14) \quad pdf(s) = \frac{\Gamma\left(\frac{L+n+1}{2}\right) exp\left\{-\frac{1}{2}(s-\gamma)!(s-\gamma)\right\}}{(2\pi)^{K_{1}/2} \pi^{n/2} \Gamma\left(\frac{L+1}{2}\right)}$$

$$\cdot \sum_{j=0}^{\infty} \frac{\left\{\frac{1}{2}(s-\gamma)!(s-\gamma)\right\}^{j}}{2j!} \int_{0}^{\infty} \frac{m^{n/2+j-1} dm}{(1+m)(L+n+1+K_{1})/2+j} \int_{V_{1,n}} (dh)$$

$$= \frac{\Gamma\left(\frac{L+n+1}{2}\right) exp\left\{-\frac{1}{2}(s-\gamma)!(s-\gamma)\right\}}{(2\pi)^{K_{1}/2} \pi^{n/2} \Gamma\left(\frac{L+1}{2}\right)}$$

$$\cdot \sum_{j=0}^{\infty} \frac{\left\{\frac{1}{2}(s-\gamma)!(s-\gamma)\right\}^{j}}{2j!} \frac{\Gamma\left(\frac{n}{2}+j\right) \Gamma\left(\frac{L+K_{1}+1}{2}\right)}{\Gamma\left(\frac{L+K_{1}+n+1}{2}+j\right)} \frac{2\pi^{n/2}}{\Gamma\left(\frac{n}{2}\right)}$$

$$= \frac{\Gamma\left(\frac{L+n+1}{2}\right) \Gamma\left(\frac{L+K_{1}+1}{2}\right) exp\left\{-\frac{1}{2}(s-\gamma)!(s-\gamma)\right\}}{(2\pi)^{K_{1}/2} \Gamma\left(\frac{L+K_{1}+1}{2}\right) \Gamma\left(\frac{n}{2}\right)} \frac{1}{1} F_{1}\left(\frac{n}{2}, \frac{n+L+K_{1}+1}{2}; \frac{1}{2}(s-\gamma)!(s-\gamma)\right).$$

$$(15)$$

It is simple to verify that the density given by (15) integrates to unity. In considering the order to which moments exist it is convenient to set $\gamma = 0$ and examine the convergence of the following series of positive terms:

$$\begin{split} E\{(s^*s)^{d/2}\} &= \frac{\Gamma\left(\frac{L+n+1}{2}\right)\Gamma\left(\frac{L+K_1+1}{2}\right)}{(2\pi)^{K_1/2}\Gamma\left(\frac{L+1}{2}\right)\Gamma\left(\frac{L+K_1+1}{2}\right)} \int_{j=0}^{\infty} \frac{\left(\frac{n}{2}\right)_{j}\left(\frac{1}{2}\right)^{j}}{j!\left(\frac{n+L+K_1+1}{2}\right)_{j}} \int_{j=0}^{c-s^*s/2} (s^*s)^{j+d/2} ds \\ &= \frac{\Gamma\left(\frac{L+n+1}{2}\right)\Gamma\left(\frac{L+K_1+1}{2}\right)}{(2\pi)^{K_1/2}\Gamma\left(\frac{L+1}{2}\right)\Gamma\left(\frac{L+K_1+n+1}{2}\right)} \int_{j=0}^{\infty} \frac{\left(\frac{n}{2}\right)_{j}\left(\frac{1}{2}\right)^{j}}{j!\left(\frac{n+L+K_1+1}{2}\right)_{j}} \frac{\Gamma\left(\frac{d+K_1}{2}+j\right)2^{K_1/2}}{2\left(\frac{1}{2}\right)} \\ &= \frac{\Gamma\left(\frac{L+n+1}{2}\right)\Gamma\left(\frac{L+K_1+1}{2}\right)\Gamma\left(\frac{d+K_1}{2}\right)2^{d/2}}{\Gamma\left(\frac{L+K_1+1}{2}\right)\Gamma\left(\frac{d+K_1+1}{2}\right)\Gamma\left(\frac{d+K_1}{2}\right)2^{d/2}} 2^{F_1\left(\frac{n}{2},\frac{d+K_1}{2},\frac{n+L+K_1+1}{2}; 1\right). \end{split}$$

The series converges absolutely provided $d \leq L$ and diverges otherwise. Thus, integer moments of γ_{IV} are finite up to the number of surplus instruments $L = K_3 - n$ (or, in the case of 2SLS, to the degree of overidentification), as we know from earlier results on β_{IV} and from the form of (5).

4. THE GENERAL CASE

From (7) and (8) the joint p.d.f. of (b, B, β_{IV}) under the alternative hypothesis $\Pi_{12} \neq 0$, $\Pi_{22} \neq 0$ is:

(16)
$$pdf(b,B,r) = \frac{\Gamma_{n}\left(\frac{L+n+1}{2}\right)etr\left\{-\frac{T}{2}(I+\beta\beta^{\dagger})\overline{\Pi}_{22}^{\dagger}\overline{\Pi}_{22}\right\}}{(2\pi)^{K_{1}(n+1)/2}\pi^{n/2}(1+r^{\dagger}r)^{(L+n+1)/2}}$$

$$exp\left\{-(b-\pi_{11})^{\dagger}(b-\pi_{11})/2\right\}etr\left\{-\frac{1}{2}(B-\Pi_{12})^{\dagger}(B-\Pi_{12})\right\}$$

$$\cdot \sum_{j=0}^{\infty} \sum_{k=0}^{\infty} \sum_{\varphi \in j \cdot \kappa} \frac{\left(\frac{L}{2}\right)_{j}\left(\frac{L+n+1}{2}\right)_{\kappa} \theta_{\varphi}^{j,\kappa}}{j!k!\Gamma_{n}\left(\frac{L+n}{2},\varphi\right)}$$

$$\cdot C_{\varphi}^{j,\kappa}\left(\frac{T}{2}\overline{\Pi}_{22}\beta\beta^{\dagger}\overline{\Pi}_{22}^{\dagger}, \frac{T}{2}\overline{\Pi}_{22}(I+\beta r^{\dagger})(I+rr^{\dagger})^{-1}(I+r\beta^{\dagger})\overline{\Pi}_{22}^{\dagger}\right).$$

As before we write γ_{IV} as s = b - Br. Using this transformation and completing the new matrix quadratic form in B we find:

The matrix B may now be integrated out leaving us with:

$$\begin{aligned} \mathrm{pdf}(s,r) &= \frac{\Gamma_{n} \left(\frac{L+n+1}{2} \right) \mathrm{etr} \left\{ -\frac{T}{2} (I+\beta\beta^{\dag}) \overline{\Pi}_{22}^{\dag} \overline{\Pi}_{22} \right\}}{(2\pi)^{K_{1}/2} \pi^{n/2} (1+r'r)} \cdot \exp \left\{ -(s-\pi_{11})^{\dag} (s-\pi_{11})/2 \right\} \mathrm{etr} \left(-\frac{1}{2} \Pi_{12}^{\dag} \Pi_{12} \right) \\ &= \mathrm{etr} \left\{ \frac{1}{2} [\Pi_{12}^{-} (s-\pi_{11}) r'] (I+rr')^{-1} [\Pi_{12}^{-} (s-\pi_{11}) r']' \right\} \\ &\cdot \sum_{j=0}^{\infty} \sum_{k=0}^{\infty} \sum_{\varphi \in j \cdot \kappa} \frac{\left(\frac{L}{2} \right)_{j} \left(\frac{L+n+1}{2} \right)_{k} \theta_{\varphi}^{j,\kappa}}{j! k! \Gamma_{n} \left(\frac{L+n}{2}, \varphi \right)} \\ &\cdot C_{\varphi}^{j,\kappa} \left(\frac{T}{2} \overline{\Pi}_{22} \beta\beta^{\dag} \overline{\Pi}_{22}^{\dag}, \frac{T}{2} \overline{\Pi}_{22} (I+\beta r') (I+rr')^{-1} (I+r\beta^{\dag}) \overline{\Pi}_{22}^{\dag} \right) \end{aligned}$$

(17)
$$= \frac{\Gamma_{n} \left(\frac{L+n+1}{2}\right) \operatorname{etr} \left\{ -\frac{T}{2} (I+\beta\beta') \overline{\Pi}_{22}^{\dagger} \overline{\Pi}_{22} \right\}}{(2\pi)^{K_{1}/2} \pi^{n/2} (1+r'r)} \operatorname{exp} \left\{ -(s-\pi_{11})'(s-\pi_{11})/2 \right\}$$

$$\cdot \operatorname{exp} \left\{ \frac{1}{2} \frac{(s-\pi_{11})'(s-\pi_{11})r'r - 2r'\Pi_{12}(s-\pi_{11}) - r'\Pi_{12}^{\dagger}\Pi_{12}r}{1+r'r} \right\}$$

$$\cdot \sum_{j=0}^{\infty} \sum_{k=0}^{\infty} \sum_{\varphi \in j \cdot \kappa} \frac{\left(\frac{L}{2}\right)_{j} \left(\frac{L+n+1}{2}\right)_{\kappa} \theta_{\varphi}^{j,\kappa}}{j!k!\Gamma_{n} \left(\frac{L+n}{2}, \varphi\right)}$$

$$\cdot C_{\varphi}^{j,\kappa} \left(\frac{T}{2}\overline{\Pi}_{22}\beta\beta'\overline{\Pi}_{22}^{\dagger}, \frac{T}{2}\overline{\Pi}_{22}(I+\beta r')(I+rr')^{-1}(I+r\beta')\overline{\Pi}_{22}^{\dagger} \right).$$

It will be convenient in what follows to use the identity:

(18)
$$(I+\beta r')(I+rr')^{-1}(I+r\beta') = I + \beta\beta' - \frac{(r-\beta)(r-\beta)'}{1+r'r}$$
.

Since the polynomial $C_{\phi}^{j,\kappa}$ is an analytic function of its matrix arguments we employ the Taylor expansion:

$$(19) \quad \operatorname{etr} \left\{ -\frac{\partial Z(\mathbf{r} - \beta)}{\partial z} (\mathbf{r} - \beta)' / (1 + \mathbf{r}' \mathbf{r}) \right\} C_{\phi}^{j, \kappa} \left(\frac{T}{2} \overline{\Pi}_{22} \beta \beta' \overline{\Pi}_{22}', \frac{T}{2} \overline{\Pi}_{22} (1 + \beta \beta' + Z) \overline{\Pi}_{22}' \right) \Big|_{Z=0}$$

$$= C_{\phi}^{j, \kappa} \left(\frac{T}{2} \overline{\Pi}_{22} \beta \beta' \overline{\Pi}_{22}', \frac{T}{2} \overline{\Pi}_{22} (1 + \beta \beta') \overline{\Pi}_{22}' - \frac{T}{2} \overline{\Pi}_{22} (\mathbf{r} - \beta) (\mathbf{r} - \beta)' \overline{\Pi}_{22} / (1 + \mathbf{r}' \mathbf{r}) \right)$$

which converges uniformly in r. The matrix Z in (19) is a matrix of auxiliary variables and ∂Z denotes the matrix operator $\partial/\partial Z$. The left side of (19) provides a simple algebraic representation of the multinomial expansion of the right hand side polynomial, which involves a sum of matrix arguments. The latter no doubt admits an expansion in terms of polynomials with more matrix arguments; but the explicit form of this expansion has not yet been derived in the multivariate literature and (19) is a simpler alternative that is very convenient for our purpose.

From (17) and (19) we obtain

$$(20) \quad pdf(s,r) = \frac{\Gamma_{n} \left(\frac{L+n+1}{2}\right) etr\left\{-\frac{T}{2}(I+\beta\beta^{\dagger}) \overline{\Pi}_{22}^{\dagger} \overline{\Pi}_{22}\right\}}{(2\pi)^{K_{1}/2} \pi^{n/2} (1+r^{\dagger}r)} etr\left\{-(s-\pi_{11})^{\dagger}(s-\pi_{11})/2\right\}$$

$$\cdot \sum_{j=0}^{\infty} \sum_{k=0}^{\infty} \sum_{\varphi \in j+\kappa} \frac{\left(\frac{L}{2}\right)_{j} \left(\frac{L+n+1}{2}\right)_{k} \theta_{\varphi}^{j,\kappa}}{j! k! \Gamma_{n} \left(\frac{L+n}{2}, \varphi\right)}$$

$$\cdot \sum_{k} \frac{\left\{\frac{1}{2}(s-\pi_{11})^{\dagger}(s-\pi_{11})\right\}^{k} \left(-\frac{1}{2}\right)^{k} (-1)^{k} 3^{+k} 4}{k_{1}! k_{2}! k_{3}! k_{4}!}$$

$$\cdot \left(\frac{r^{\dagger}r}{1+r^{\dagger}r}\right)^{k} \left(\frac{r^{\dagger} \overline{\Pi}_{12}^{\dagger} \overline{\Pi}_{12}^{r}}{1+r^{\dagger}r}\right)^{k} \left(\frac{r^{\dagger} \overline{\Pi}_{12}^{\dagger} (s-\pi_{11})}{1+r^{\dagger}r}\right)^{k} \left(\frac{(r-\beta)^{\dagger} \partial Z(r-\beta)}{1+r^{\dagger}r}\right)^{k} 4}$$

$$\cdot C_{\varphi}^{j,\kappa} \left(\frac{T}{2} \overline{\Pi}_{22}^{\dagger} \beta\beta^{\dagger} \overline{\Pi}_{22}^{\dagger}, \frac{T}{2} \overline{\Pi}_{22}^{\dagger} (I+\beta\beta^{\dagger}+Z) \overline{\Pi}_{22}^{\dagger}\right) \Big|_{Z=0}$$

where \sum_{ℓ} denotes $\sum_{\ell=1}^{\ell}$, $\ell=1$. Since the series converges uniformly in r we may integrate termize to remove r. The typical term is then:

(21)
$$\int \frac{(\mathbf{r'}\Pi_{12}^{\dagger}\Pi_{12}\mathbf{r})^{\ell_{2}}\{\mathbf{r'}\Pi_{12}(\mathbf{s}-\pi_{11})\}^{\ell_{3}}\{(\mathbf{r}-\beta)^{\dagger}\partial\mathbf{Z}(\mathbf{r}-\beta)\}^{\ell_{4}}(\mathbf{r'}\mathbf{r})^{\ell_{1}}d\mathbf{r}}{(\mathbf{l}+\mathbf{r'}\mathbf{r})^{(L+n+K_{1}+1)/2+\ell}}$$

$$= (-\partial\mathbf{x'}\Pi_{12}^{\dagger}\Pi_{12}\partial\mathbf{x})^{\ell_{2}}\{-\mathbf{i}\partial\mathbf{x'}\Pi_{12}(\mathbf{s}-\pi_{11})\}^{\ell_{3}}\{(\mathbf{i}\partial\mathbf{x}+\beta)^{\dagger}\partial\mathbf{Z}(\mathbf{i}\partial\mathbf{x}+\beta)\}^{\ell_{4}}$$

$$\cdot \int \frac{e^{\mathbf{i}\mathbf{x'}\mathbf{r}}(\mathbf{r'}\mathbf{r})^{\ell_{1}}d\mathbf{r}}{(\mathbf{l}+\mathbf{r}+\mathbf{K_{1}}+1)/2+\ell}\Big|_{\mathbf{x}=0}$$

where $\ell=\ell_1+\ell_2+\ell_3+\ell_4$ and ∂x denotes the operator $\partial/\partial x$ taken with respect to a vector of auxiliary variables x.

We transform $r \rightarrow Hr = p$ for H orthogonal and integrate over

the orthogonal group O(n), normalized so that the measure over the whole group is unity. The latter measure will be denoted by (\underline{dH}) . (21) becomes

$$\int_{p} \left[(-\partial x' \Pi_{12}^{\prime} \Pi_{12} \partial x)^{\ell_{2}} \{ -i \partial x' \Pi_{12} (s - \pi_{11}) \}^{\ell_{3}} \right]$$

$$\cdot \{ (i \partial x + \beta)^{\prime} \partial Z (i \partial x + \beta) \}^{\ell_{4}} \int_{O(n)} etr(iHxp') (\underline{dH}) \Big]_{x=0} \frac{(p'p)^{\ell_{1}} dp}{(1 + p'p)^{(L + n + K_{1} + 1)/2 + \ell}}$$

$$= \int_{p} \left[(-\partial x' \Pi_{12}^{\prime} \Pi_{12} \partial x)^{\ell_{2}} \{ -i \partial x' \Pi_{12} (s - \pi_{11}) \}^{\ell_{3}} \right]$$

$$\cdot \{ (i \partial x + \beta)^{\prime} \partial Z (i \partial x + \beta) \}^{\ell_{4}} \int_{0} F_{1} \left(\frac{n}{2}; -\frac{1}{4} x' x p' p \right) \Big]_{x=0} \frac{(p'p)^{\ell_{1}} dp}{(1 + p'p)^{(L + n + K_{1} + 1)/2 + \ell}}$$

$$(22) = \sum_{\tau} \left[(-\partial x' \Pi_{12}^{\prime} \Pi_{12} \partial x)^{\ell_{2}} \} -i \partial x' \Pi_{12} (s - \pi_{11}) \right]^{\ell_{3}}$$

$$\cdot \{ (i \partial x + \beta)^{\prime} \partial Z (i \partial x + \beta) \}^{\ell_{4}} \left(-\frac{1}{4} x' x \right)^{t} \Big]_{x=0} \frac{1}{t! \left(\frac{n}{2} \right)_{+}} \int_{p} \frac{(p'p)^{\ell_{1} + t} dp}{(1 + p'p)^{(L + n + K_{1} + 1)/2 + \ell}}$$

where the summation $\sum_{\mathbf{T}}$ is over all values of t for which the quantity in square brackets is non zero. Since the latter quantity is zero whenever $t > \ell_2 + \ell_3 + \ell_4$ the integral in (22) is convergent within this summation. Upon evaluation of this integral (as in (14) above) we obtain:

$$(23) \quad \frac{\pi^{n/2}}{\Gamma\left(\frac{n}{2}\right)} \sum_{\tau} \left[\left(-\partial x^{\tau} \Pi_{12}^{\tau} \Pi_{12} \partial x \right)^{\ell_{2}} \left\{ -i \partial x^{\tau} \Pi_{12} \left(s - \pi_{11} \right) \right\}^{\ell_{3}} \right. \\ \left. \cdot \left\{ \left(i \partial x + \beta \right)^{\tau} \partial Z \left(i \partial x + \beta \right) \right\}^{\ell_{4}} \left(-\frac{1}{4} x^{\tau} x \right)^{t} \right]_{x=0} \frac{1}{t! \left(\frac{n}{2}\right)_{t}} \\ \cdot \frac{\Gamma\left(\ell_{1} + \frac{n}{2} + \ell_{1}\right) \Gamma\left(\frac{L + K_{1} + 1}{2} + \ell_{2} + \ell_{3} + \ell_{4} - t\right)}{\Gamma\left(\frac{L + n + K_{1} + 1}{2} + \ell_{1} + \ell_{2} + \ell_{3} + \ell_{4}\right)} .$$

From (20) and (23) we deduce the following general expression for the p.d.f. of $\gamma_{\rm IV}$:

$$\begin{aligned} \text{(24)} \quad & \text{pdf(s)} \, = \, \frac{\Gamma_{\text{n}} \! \left(\frac{L + \text{n} + 1}{2} \right) \text{etr} \! \left\{ -\frac{T}{2} \! \left(1 + \beta \beta^{\, \, i} \right) \overline{\Pi}_{22}^{\, \, i} \overline{\Pi}_{22}^{\, \, i} \right\}}{\left(2 \pi \right)^{K_{1}/2} \Gamma \! \left(\frac{n}{2} \right)} \, \exp \left\{ - \left(s - \pi_{11} \right)^{\, \, i} \left(s - \pi_{11} \right)^{\, \, \, i} \left(s - \pi_{11} \right)^{\, \, i} \left(s - \pi_{11} \right$$

The leading case that occurs when $\Pi_{12}=0$, $\Pi_{22}=0$ may be deduced from (24) by noting that non zero terms in the various summations arise only when $j=k=\ell_2=\ell_3=\ell_4=t=0$. Moreover, $\Gamma_n((L+n)/2;\,\phi)=\Gamma_n((L+n)/2)$ when ϕ is a partition of zero. We find in this case

which is the same as (15) derived earlier by direct methods.

REMARKS

The exact densities (15) and (24) relate to the standardized model.

The corresponding densities for the non standardized model may be obtained from these results by transformation using the formulae in Phillips (1983c).

Accurate approximations to these densities that will permit wide ranging numerical computations and the analysis of marginal distributions are the next step in studying these distributions. Methods used by the author (1983d) elsewhere seem promising in this respect and will be the subject of further work.

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